

Quasi 9 and 30–40 days periodicities in the solar differential rotation

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Abstract

Using the daily Mt. Wilson Doppler velocity data during 1986–1994 (solar cycle 22), we studied the short-term variations of the order of a few days to a month timescales in the solar differential rotation. We represent the differential rotation in the form: $\omega(\lambda) = \bar{A} + \bar{B}(5 \sin^2 \lambda - 1) + \bar{C}(21 \sin^4 \lambda - 14 \sin^2 \lambda + 1)$, using a set of Gegenbauer polynomials, where $\omega(\lambda)$ is the angular velocity at latitude λ . The coefficients \bar{A} , \bar{B} , and \bar{C} are free of crosstalk. We found that ≈ 9 -day periodicity is statistically highly significant in the variations of \bar{C} at the maximum of solar cycle 22. A similar periodicity is found in the variations of \bar{B} during the descending phase of the cycle 22 with significant on $\geq 99.9\%$ confidence level. At this cycle maximum, a 30–40 day periodicity is found to be dominant among the variations in \bar{B} , and this periodicity is found in \bar{A} during almost throughout the period 1986–1994. The ≈ 9 -day periodicity in the variation of the differential rotation approximately matches with the known quasi 10-day periodicity in the total solar irradiance (TSI) variability. Hence, we speculate that there exists a relationship between the differential rotation and TSI variability. We suggest that the 9–10 day periodicities of the differential rotation and TSI have a relationship with the production and the emergence rates of the large-scale solar magnetic flux.

Keywords: solar rotation, solar magnetic field, solar activity

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1. Introduction

Besides the well known ≈ 11 -year solar cycle, solar activity varies on many short-time scales (a few minutes to a few years) and many long time scales (a few decades to the Sun's evolution time scale). It has been generally accepted that interactions of the Sun's differential rotation (DR) and the solar magnetic field play a basic role in generation of solar activity. However, the exact role of the DR in the variations of the solar activity is not known and the mechanism of solar cycle is not yet fully understood. Therefore, studies on the temporal variations of the DR are important for understanding the role of the DR in the solar variability. So far many attempts have been made to determine the variations in the DR and detected several periodicities in the DR by using the different data and different techniques (*e.g.*, Howard & LaBonte, 1980; LaBonte & Howard, 1982; Gilman & Howard, 1984; Balthasar, Vázquez, & Wöhl, 1986; Komm, Howard, & Harvey, 1993; Javaraiah & Gokhale, 1995; Javaraiah & Komm, 1999; Gupta, Sivaraman, & Howard, 1999; Howe et al., 2000a; Javaraiah, 2003, 2005; Javaraiah, Bertello, & Ulrich, 2005; Brajša, Ruždjak, & Wöhl, 2006; Javaraiah & Ulrich, 2006; Tlatov, 2007; Javaraiah et al., 2009; Chandra, Vats, & Iyer, 2010). A higher rotation rate during the cycle minimum than during the maximum is found in many studies (*e.g.*, Brajša, Ruždjak, & Wöhl, 2006).

Variation of the DR with the 11-year solar cycle is well established now. The 11-year torsional oscillations were discovered by Howard & LaBonte (1980), using Mt. Wilson Doppler velocity measurements have been confirmed using different data sets and methods (see Javaraiah & Gokhale, 2002). The torsional oscillations consist of alternating bands of faster and slower than average rotation moving from high latitudes toward the equator. The faster than average bands start at mid-latitudes of about 45° during the minimum phase of the solar cycle and move equatorward with the rising phase of the new cycle. The bands extend over about 10° in latitude. The faster than average rotation band is located on the equatorward side of the magnetic activity belt and a slower than average rotation band is located on its poleward side. Thus, it is thought that the torsional oscillations are associated with magnetic shear. At high latitudes, the rotation changes from slower to faster than average at (or just after) the solar maximum, which might be related to the polarity reversal of polar magnetic fields, and remains faster than average to the next solar minimum without clear migratory character. Helioseismic observations show that the tor-

sional oscillations are not just a superficial phenomenon but they extend to at least the upper third of the solar convection zone (Howe et al., 2000b; Antia, Basu, & Chitre, 2008). Variations on the few other time scales in the coefficients of solar differential rotation (Javaraiah & Gokhale, 1995, 1997; Javaraiah & Komm, 1999; Javaraiah, Bertello, & Ulrich, 2005) and the residual rotation (Brajša, Ruždjak, & Wöhl, 2006), including one which is approximately equal to the Gleissberg cycle, have been found using sunspot group data.

It is known that the total solar irradiance (TSI) varies by about 0.1% over the solar cycle. A number of statistical models of TSI variability have been constructed on the basis of inhomogeneities of surface magnetic field. These models helped to identify the surface magnetic structures (sunspots and faculae) responsible for variations in TSI and also provided widely accepted theoretical explanations. However, the exact mechanism behind the TSI variation is not known (Foukal, 1992; Kuhn, Lin, & Coulter, 1999; Solanki, Krivova, & Wenzler, 2005; Domingo et al., 2009).

Kuhn, Libbrecht, & Dicke (1988) found that a significant part of the solar cycle variation in TSI is a result of the temporal changes in the latitude-dependent surface temperature of the Sun. These authors speculated that the DR (turbulent Reynolds stresses and their perturbations) has a role in TSI variability. Besides the ≈ 11 -year periodicity, TSI also seems to vary on timescales which are much shorter than 11-year. For example, a ≈ 10 -day periodicity is found to be prominent in the short variations of TSI (Foukal, 1992; Nikonova, Klochek, & Palamrchuk, 1998) and the existence of a ≈ 30 -day periodicity in the proxies of TSI is reported (*e.g.*, Pap, Willson, & Donnelly, 1992). However, the amplitudes of the short variations of TSI depend on the phase of a solar cycle (*e.g.*, Pap, Willson, & Donnelly, 1992). In the present analysis we have attempted to find the short periodicities in the coefficients of the DR determined from the Mt. Wilson velocity data during 1986–1994, because which may help for understanding the TSI variability.

In the next section we describe the methodology and the data analysis. In Section 3 we present the results, and in Section 4 we present conclusions and briefly discuss them.

2. Methodology and data analysis

The Sun’s DR can be determined from full disc velocity data using the traditional polynomial expansion:

$$\omega(\lambda) = A + B \sin^2 \lambda + C \sin^4 \lambda, \quad (1)$$

where $\omega(\lambda)$ is the solar rotation rate at latitude λ , the parameter A represents the equatorial rotation rate, B and C measure the latitude gradient of the rotation rate with B representing mainly low latitudes and C representing largely higher latitudes.

As pointed out by several authors (see Snodgrass, 1984), due to the non-orthogonality of the fit functions, the coefficients A , B , and C have a crosstalk which affects their temporal behavior. Snodgrass & Howard (1985) used the so-called Gegenbauer Polynomials (Morse & Feshbach, 1953) as a set of disk-orthogonal fit functions with $T_0^1(\sin \lambda) = 1$, $T_2^1(\sin \lambda) = 5 \sin^2 \lambda - 1$, and $T_4^1(\sin \lambda) = 21 \sin^4 \lambda - 14 \sin^2 \lambda + 1$, which leads to the following expansion:

$$\omega(\lambda) = \bar{A} + \bar{B}(5 \sin^2 \lambda - 1) + \bar{C}(21 \sin^4 \lambda - 14 \sin^2 \lambda + 1). \quad (2)$$

The coefficients \bar{A} , \bar{B} , and \bar{C} are free of crosstalk, \bar{A} represents the ‘rigid body’ (or ‘mean’) component in the rotation, \bar{B} and \bar{C} are the components of the DR. If the polynomial expansion is terminated at \bar{C} (or C), the coefficients, \bar{A} , \bar{B} , and \bar{C} are related to the A , B , and C coefficients as follows:

$$\bar{A} = A + (1/5)B + (3/35)C; \quad \bar{B} = (1/5)B + (2/15)C; \quad \bar{C} = (1/21)C. \quad (3)$$

In this case, the temporal variation of \bar{C} is qualitatively identical to that of C (Note: in Equation (6) of Snodgrass & Ulrich (1990) $1/21$ is erroneously typed as $2/21$).

Earlier, using the daily values of the A , B , C coefficients (cf., Equation (1)) derived from the Mt. Wilson velocity data during the period 1967–1994, Javaraiah & Komm (1999) determined nearly one-year and more than one year periodicities in the DR. In that early paper the authors had focused mainly on the data obtained during the period 1982–1994, after a major change in the Mt. Wilson spectrograph modification which reduced the instrumental noise (see Howard et al., 1983). Recently, Javaraiah et al. (2009) found the existence of a ≈ 1.4 -year periodicity in A determined from the Mt. Wilson velocity data during 1986–1995, but they did not find any significant variation in A determined from the data during 1996–2007, which are

measured from the more stable Mt. Wilson spectrograph instrumentation. However, the 1.3-year periodicity in the variation of the low and the middle latitudes' rotation rate near the base of the convection zone, that was detected by Howe et al. (2000a) from helioseismic measurements during the period 1995–1999, is also found to be not persisting after 2001 (Howe et al., 2007). In fact, in several solar activity phenomena a 1.3-year periodicity has been found to be dominant during the cycle 22 and weak or absent in the later period (see Obridko & Shelting, 2007).

The time series of A , B , and C during 1986–1994 have many data gaps that vary in size, 1–18 days. Since here we have attempted to determine the periodicities of the order of a few days to a month only, hence it is necessary to analyse the daily data which have no gaps or have only a few 1–2 day gaps. We find that the data in the following time (rotation number) intervals contain only a few (two or three) 1–2 days gaps: 7772–7934, 8197–8297, 8559–8726, 8902–9023, 9280–9376, 9583–9720, 10056–10161, 10375–10514 and 10731–10893. The size of each of these intervals is sufficient to determine the short periodicities of a few days to a month time. (In the data during a period which is not included in any one of these intervals there are large number of gaps of different sizes. Therefore, they are unfit to use for the present purpose.) Incidentally, each one of these intervals belongs to one of the years 1986–1994, in chronological order. That is, the years 1986, 1987,, 1994 contain 7772–7934 (first interval), 8197–8297 (second interval), ..., 10731–10893 (last interval), respectively. By using Equation (3), we converted the A , B , and C coefficients to the corresponding coefficients \bar{A} , \bar{B} , and \bar{C} , respectively. We corrected the time series of each of these coefficients by removing the very large-spikes, i.e., we have removed the values which are $> 2.5\sigma$ level (where σ is the standard deviation of the original time series). We filled the gaps in each of these time-series from the values obtained by liner-interpolations. Figures 1 and 2 show these corrected time series and also the corresponding uncorrected time-series. As can be seen in these figures, there are variations of the order of few days to a month time scales in \bar{A} , \bar{B} and \bar{C} during many of the aforementioned time intervals. We computed fast Fourier transform (FFT) power spectra of \bar{A} , \bar{B} , and \bar{C} . The results are presented in the next section.

3. Results

Figures 3–5 show the FFT power spectra of \bar{A} , \bar{B} and \bar{C} derived from the corrected data (solid curves in Figures 1–2) during each of the reliable time intervals found above. Before computing the FFT, the long-term trend of the order of the length of the time series was removed by subtracting the corresponding leaner-model of the series and a cosine bell function was applied to the first and the last 10% of the time series (see Brault & White, 1971). In Table 1 we have given the periodicities in the variations of \bar{A} , \bar{B} and \bar{C} , whose levels of significance are $\geq 2\sigma$ in the respective FFT power spectra. The results in this table (and Figures 3–5) suggest that at the solar cycle maximum (*i.e.*, in the interval 8902–9023 during the year 1989), a ≈ 9 -day periodicity is very pronounced in the variations of \bar{C} . In the same time this periodicity is insignificant/absent in the variations of both \bar{A} and \bar{B} , and a 30–40 day periodicity is dominant in the variations of \bar{B} . During the decay phase of the cycle 22 (years 1991–1993), there is ~ 9 -day periodicity in the variations of \bar{B} with $> 99\%$ confidence level. There is a suggestion on the existence of the 30–40 day periodicity in \bar{A} during almost throughout the period 1986–1994.

4. Conclusions and discussion

By analyzing the Mt. Wilson Doppler velocity data during the period 1986–1994, we find:

1. During the maximum phase of the solar cycle 22 there were highly statistically significant ≈ 9 -day and 30–40 day periodicities in the DR coefficients \bar{C} and \bar{B} , respectively.
2. During the descending phase of cycle 22 the ≈ 9 -day periodicity was stronger in \bar{B} than in \bar{C} , the 30–40 day periodicity was weak in both \bar{B} and \bar{C} .
3. Near the minimum of cycle 22 the 30–40 day periodicity was present and the ≈ 9 -day periodicity was insignificant/absent in \bar{C} .
4. In \bar{A} the 9-day periodicity was absent during any phase of the solar cycle 22, whereas a 30–40 day periodicity was present in almost throughout this cycle.

It is difficult to detect variations in the equatorial rotation rate determined from the Doppler velocity measurements because several observational

and instrumental effects can produce spurious peaks with similar periodicities (*e.g.*, Snodgrass, 1983; Lustig & Wöehl, 1989; Snodgrass & Ulrich, 1990; Stenflo, 1990; Ulrich & Bertello, 1996; Wöhl & Schmidt, 200). Variations in the solar rotation on shorter time scales are more difficult to detect because of the relatively low amplitudes and intermittent nature of them.

Since 1986 the quality of the Mt. Wilson velocity data is high enough to study substantially short variations in the derived velocity values (Snodgrass, 1992; Ulrich, 1998), hence, the dominant ≈ 9 day and the 30–40 day periodicities in \bar{B} and \bar{C} , found above, may be really exist in the DR. However, the amplitudes of these periodicities in the DR depend on the phases of the solar cycle.

A ≈ 10 -day periodicity seems to prominently present in the variation of the TSI (see Foukal, 1992; Nikonova, Klochek, & Palamrchuk, 1998). Pap, Willson, & Donnelly (1992) found a ≈ 30 -day periodicity in the proxies of TSI. Recently, Shapiro et al. (2011) found 13- and 27-day periodicities in the spectral solar irradiance. The ≈ 9 -day and 30–40 day periodicities of \bar{B} and \bar{C} approximately match with that known ≈ 10 -day and ≈ 30 -day periodicities in the TSI variability.

The existence of the common periodicities in the DR and the TSI suggests that the latter may be related to the former as pointed by Kuhn, Libbrecht, & Dicke (1988). Hence, the ≈ 9 -day and 30–40 day periodicities in the DR, which are detected here, may be of particular importance for understanding the TSI variability.

It may be also necessary to note here that the ≈ 9 day⁻¹ frequency in the DR may be the third harmonic of the Sun’s rotation frequency. The ≈ 9 -day and ≈ 30 -day periodicities seem to be predominantly present in the solar coronal holes, solar wind, auroral electron power, and geomagnetic parameters (Temmer, Vršnak, & Veronig, 2007; Emery et al., 2009). Also, helioseismic studies suggest that global solar oscillation frequency changes on time scale as short as nine days (Tripathy et al., 2007).

Howard & LaBonte (1981) analysed Mt. Wilson magnetograph data during the period 1967 to mid-1980 and found that the rate at which the magnetic flux appears on the Sun is sufficient to create all the flux that is seen at the solar surface within a period of about 10 days. The magnetic structures of the sunspot groups may rise from near the bottom of the convection zone to the surface in about 10 days (Javaraiah & Gokhale, 1997). Thus, the 9–10 day periodicity of the DR and the TSI seems to be related to the production and emergence rate of the large-scale solar magnetic flux.

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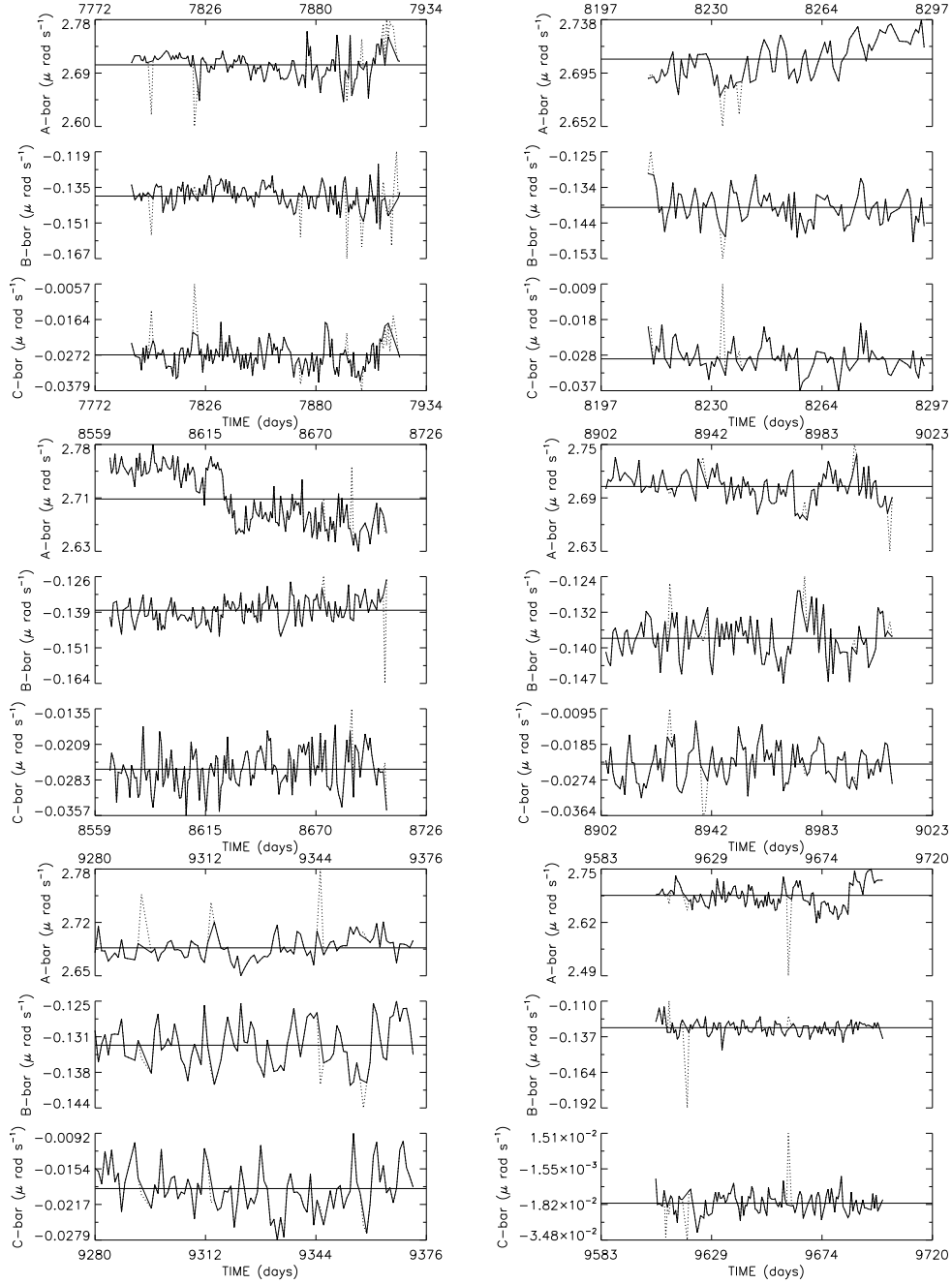


Figure 1: Variations in the \bar{A} , \bar{B} and \bar{C} (cf. Equation (3)) determined from Mt. Wilson velocity data in the time (rotation number) intervals 7772–7934 (top panel-left), 8197–8297 (top panel-right), 8559–8726 (middle panel-left), 8902–9023 (middle panel-right), 9280–9376 (bottom panel-left), and 9583–9720 (bottom panel-right). The dotted- and solid-curves represent the variations determined from the uncorrected and the corrected data, respectively. The horizontal lines represent the mean values of the corrected data.

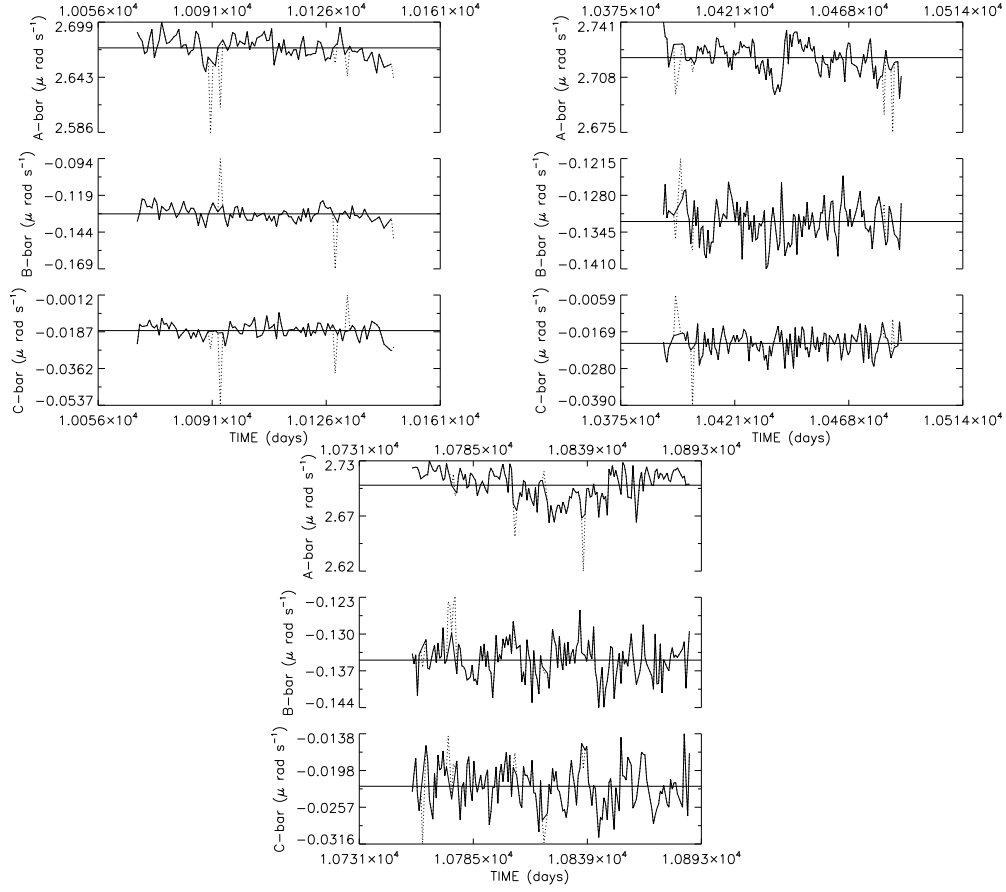


Figure 2: The same as Figure 1, but for the intervals 10056–10161 (top panel-left), 10375–10514 (top-panel-right), and 10731–10893 (bottom panel).

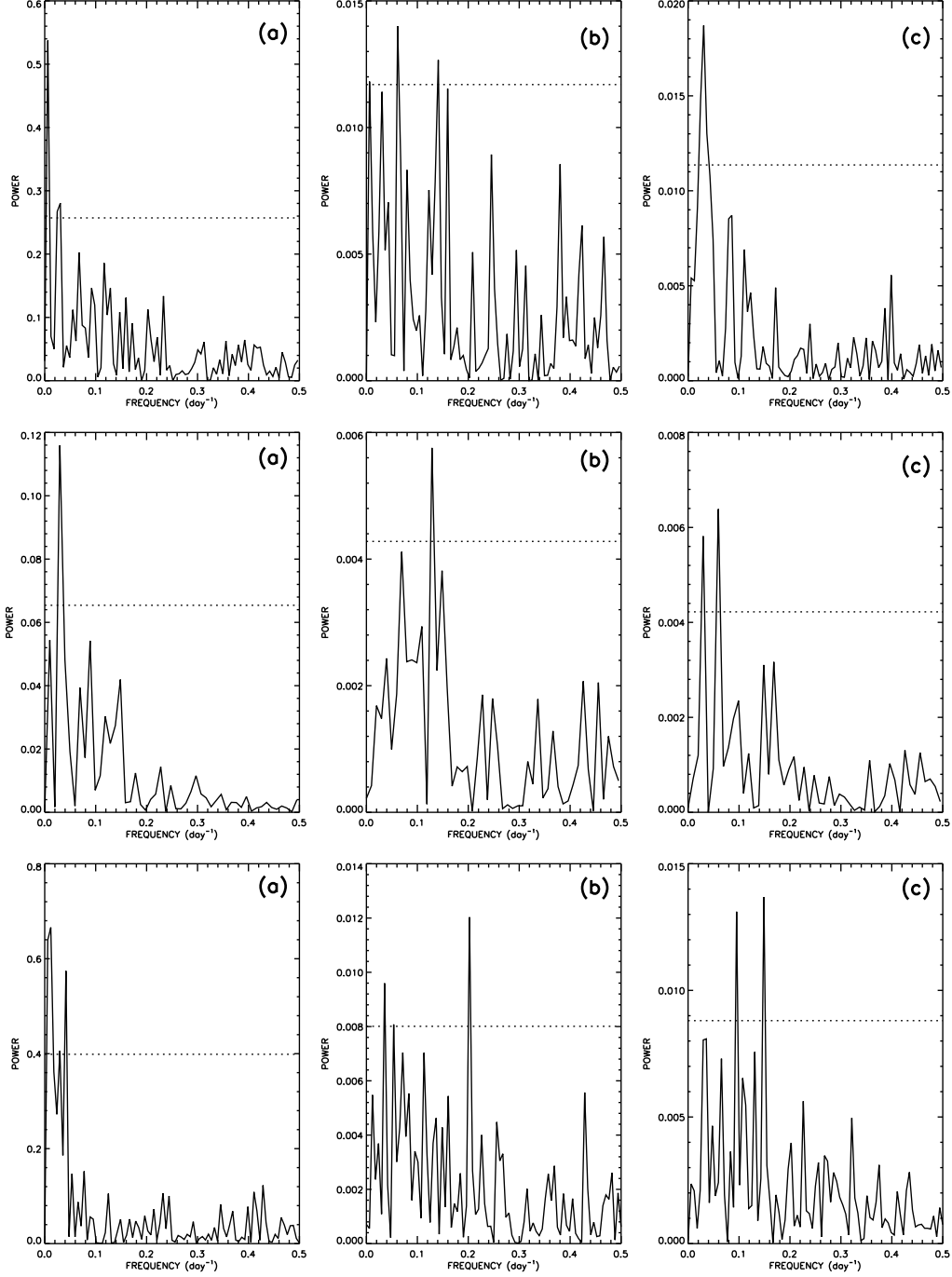


Figure 3: FFT power spectra of (a) \bar{A} , (b) \bar{B} and (c) \bar{C} determined from the Mt. Wilson velocity data in the rotation number intervals 7772–7934 (top panel), 8197–8297 (middle panel), and 8559–8726 (bottom panel). The dotted horizontal line represents 99% confidence level.

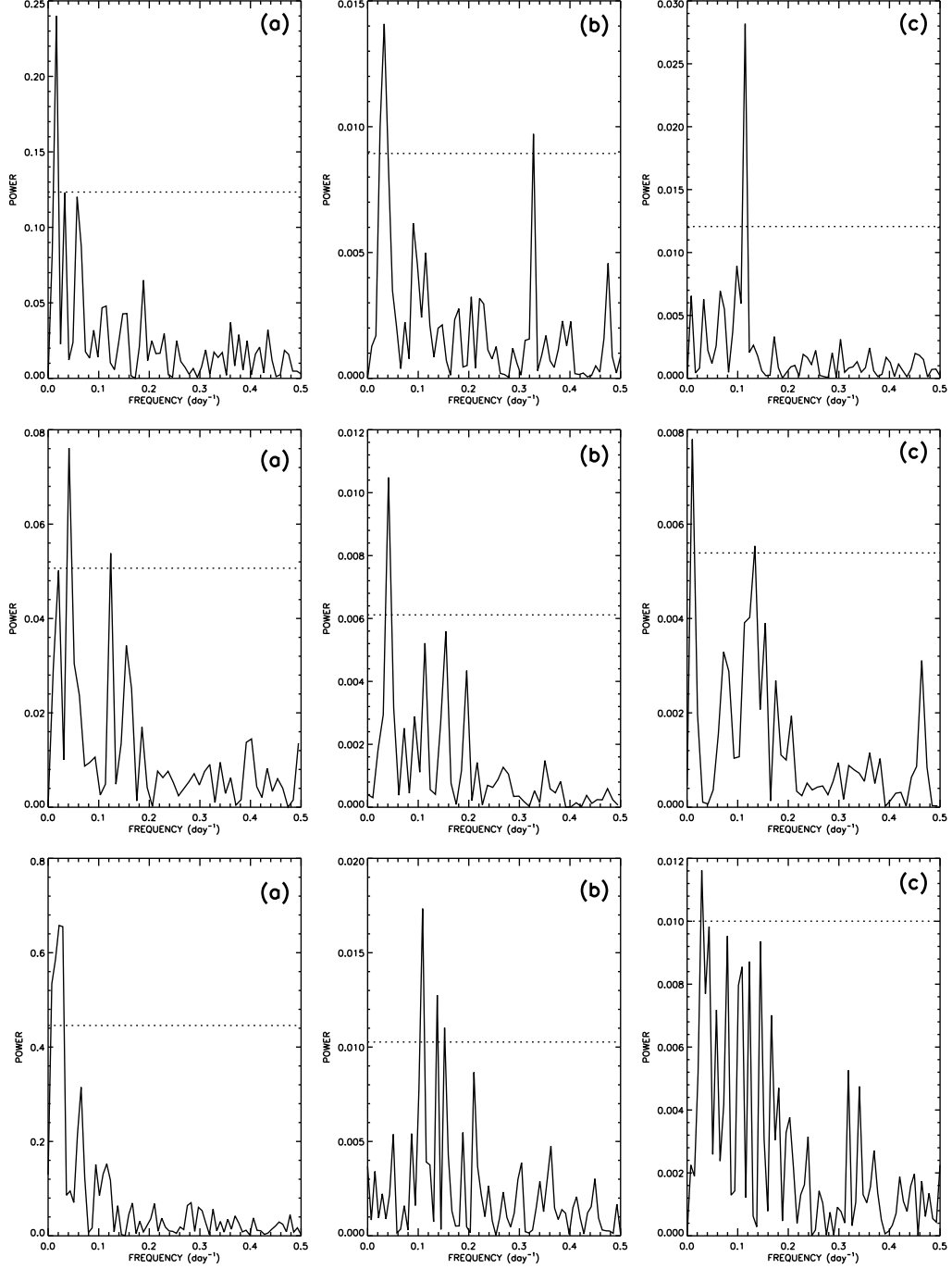


Figure 4: The same as Figure 3, but for the intervals 8902–9023 (top panel), 9280–9376 (middle panel), and 9583–9720 (bottom panel).

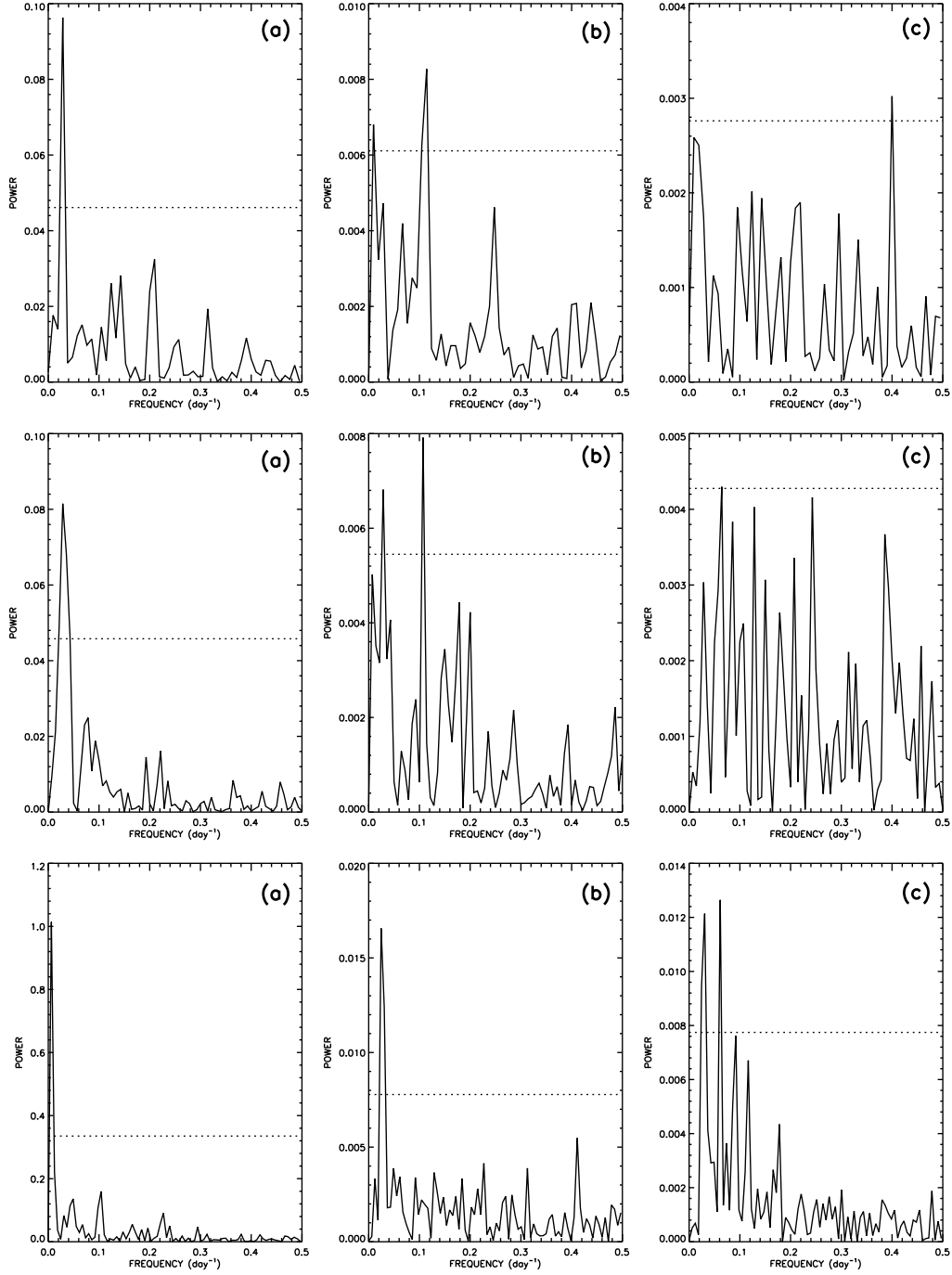


Figure 5: The same as Figure 3, but for the intervals 10056–10161 (top panel), 10375–10514 (middle panel), and 10731–10893 (bottom panel).

Table 1: Periodicities in the variations of \bar{A} , \bar{B} , and \bar{C} determined from the daily Mt. Wilson velocity data in different time intervals mentioned in the first column (the corresponding year and number of data points are mentioned within round and square brackets, respectively). The levels of significance (in units of σ) are also given.

Time Interval	\bar{A} Period (in days)	σ	\bar{B} Period (in days)	σ	\bar{C} Period (in days)	σ
7772–7934 (1986), [163]	~ 163 33–41	6.2 2.7–2.9	~ 163 $\sim 33 \pm 8$ ~ 16 6–7	3.7 2.4 3.3 2.5–2.9	33–41 23–27	3.5–4.7 2.4–3.0
8197–8297 (1987), [101]	~ 101 ~ 34 ~ 11	2.0 2.0 2.0	~ 14 6–7	2.4 2.2–3.84	~ 34 ~ 17	3.8 4.3
8559–8726 (1988), [168]	56–84 ~ 34 ~ 24	2.3–4.6 2.6 3.9	~ 28 ~ 19 ~ 9 ~ 5	3.3 2.7 2.1 4.3	28–34 ~ 10 6–8	2.9 4.3 2.0–4.5
8902–9023 (1989), [122]	~ 61 ~ 30 ~ 17	5.6 2.6 2.5	~ 30 ~ 24 ~ 3	4.5 2.4 2.9	~ 9	6.7
9280–9376 (1990), [97]	~ 48 ~ 24 ~ 8	2.5 4.3 2.8	~ 24 ~ 9 ~ 6	4.9 2.1 2.3	~ 97 ~ 7	4.1 2.7
9583–9720 (1991), [138]	34–69	3.5–4.0	~ 9 6–7 ~ 4	4.9 2.8–3.3 2.0	~ 34 ~ 12 8–9 ~ 6	3.1 2.4 2.0 2.3
10056–10161 (1992), [106]	~ 35	5.0	~ 35 8–9	3.4 2.7–3.8	52–105 ~ 2.5	2.2–2.3 2.9
10375–10514 (1993), [140]	~ 35	5.0	~ 140 ~ 35 ~ 9	2.3 3.4 4.1	~ 15 ~ 12 ~ 8 ~ 4 ~ 2.5	2.6 2.2 2.4 2.5 2.0
10731–10893 (1994), [163]	~ 163	8.4	~ 32 –41	4.5–6.3	~ 32 –41 ~ 16 ~ 11 ~ 9	3.3–4.4 4.6 2.5 2.1